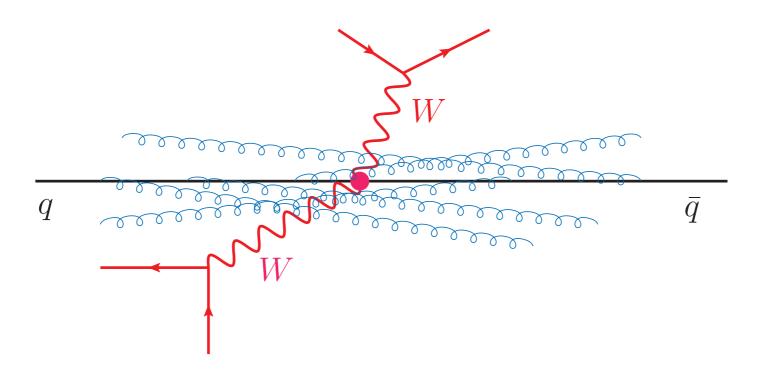
Resummation Predictions for WWW Production at the LHC

Prerit Jaiswal Brown University

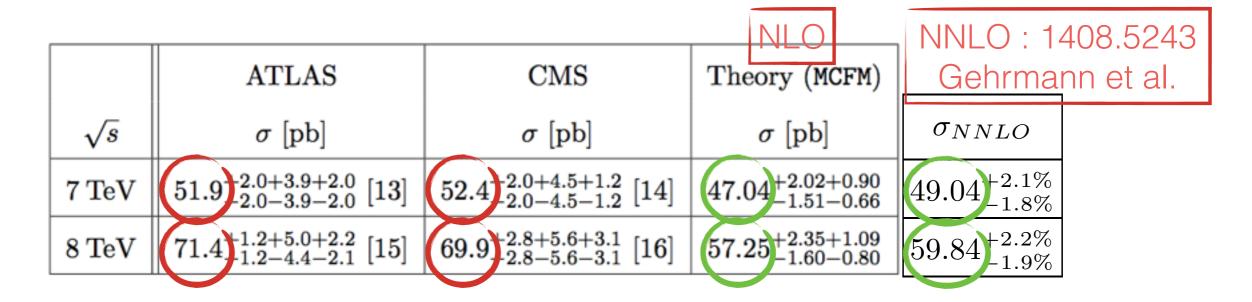


References:

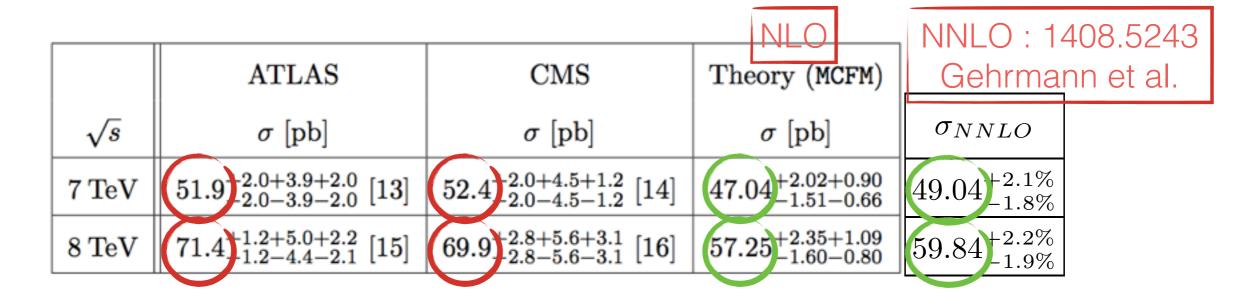
arXiv:1407.4537 : PJ and Takemichi Okui

arXiv:1506.07529 : PJ and Takemichi Okui

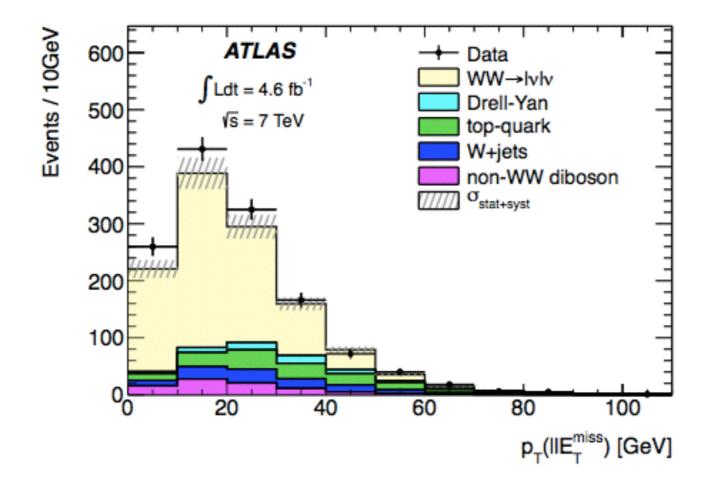
arXiv:1509.07118: PJ, Patrick Meade and Harikrishnan Ramani



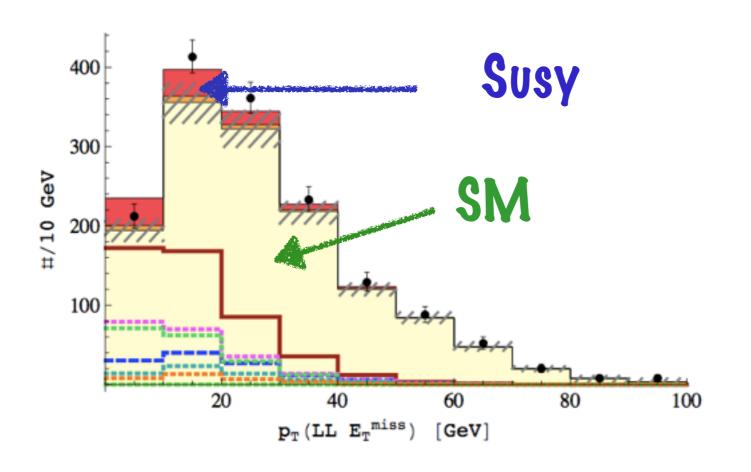
- Process : pp → WW → ℓ v ℓ v
- Mild excesses reported by ATLAS <u>and</u> CMS at 7 <u>and</u> 8 TeV measurements.
- Discrepancy reduces slightly at NNLO but does not go away.

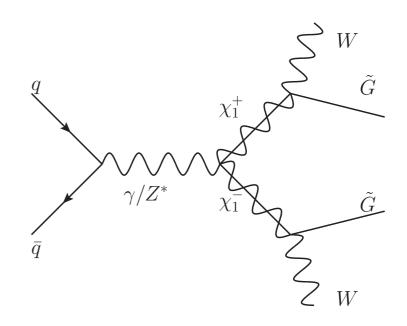


Discrepancy also exists in shapes.



- New physics hiding in plain sight? (ℓ ℓ + MET final state)
- Could it be SUSY?



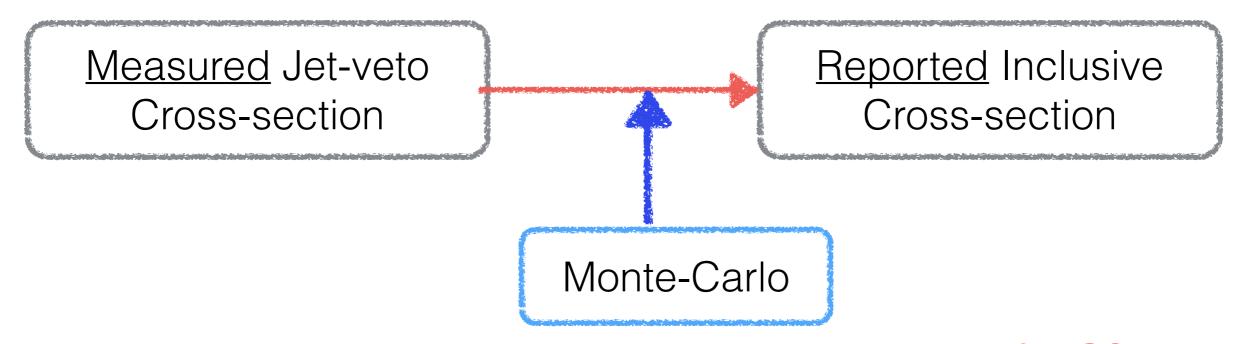


110 GeV charginos

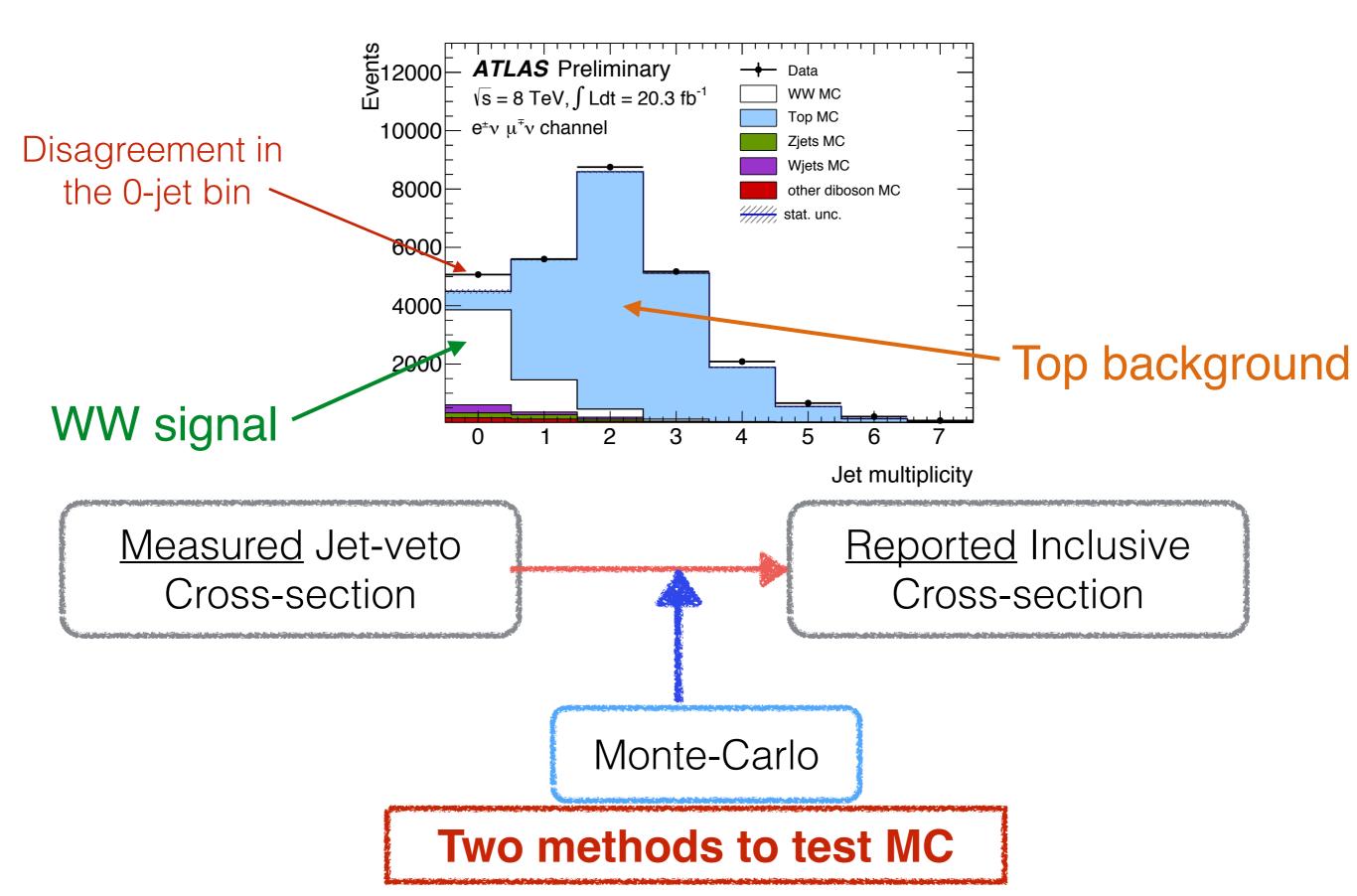
D. Curtin, PJ, and P. Meade, Charginos hiding in plain sight [arXiv:1206.6888]

 Any new physics charged under electroweak gauge group could possibly lead to such signatures. Other proposed explanations for the WW excess include sleptons and stops.

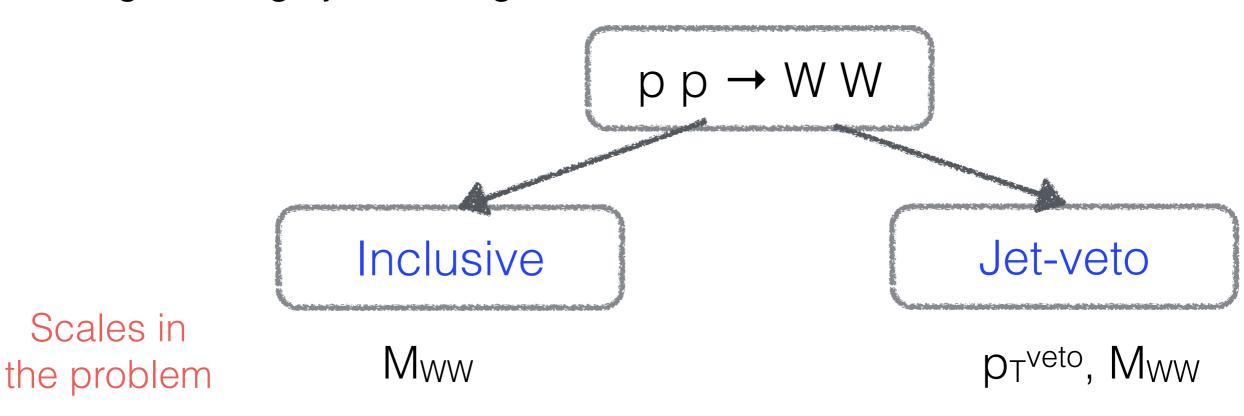
- Could the WW excess have a QCD explanation?
- Cross-section reported : p p → W W + X
 X are all hadronic final states i.e. inclusive measurement
- Actual measurement : p p → W W + X'
 X' are <u>some</u> hadronic final states that pass jet-veto condition.



Do we have a good theoretical understanding of MC?



Origin of large jet veto logs :



Logs at

higher orders

$$\log(M_{\it WW}/\mu)$$

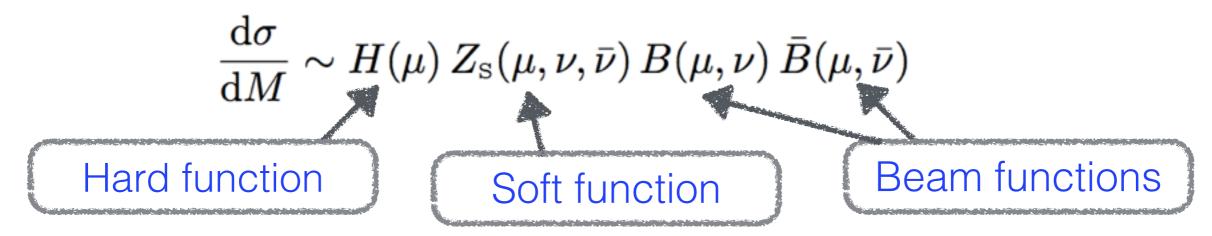
$$\log(M_{WW}/\mu)$$
$$\log(p_T^{veto}/\mu)$$

Choice of μ to minimize logs

$$\mu \sim M_{WW}$$

No choice of μ Large logs of the form $\log(p_T^{veto}/M_{ww})$ remain

- Our approach: Calculate jet-veto cross-section analytically by resummation at NNLL using SCET without relying on MC. [arXiv:1407.4537: PJ and T. Okui]
- Factorization of cross-section :

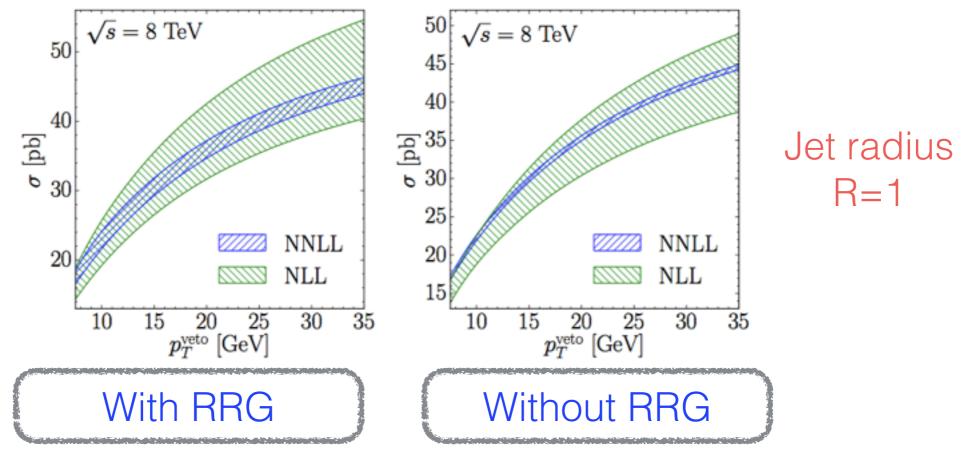


- Beam functions have divergences (rapidity divergences) which are <u>not</u> regulated by dimensional regularization
- We introduce the following analytic regulator

$$\left(\frac{
u}{k_+}\right)^{\!\!lpha} heta(k_+ - k_-) + \left(\frac{ar
u}{k_-}\right)^{\!\!arlpha} heta(k_- - k_+)$$

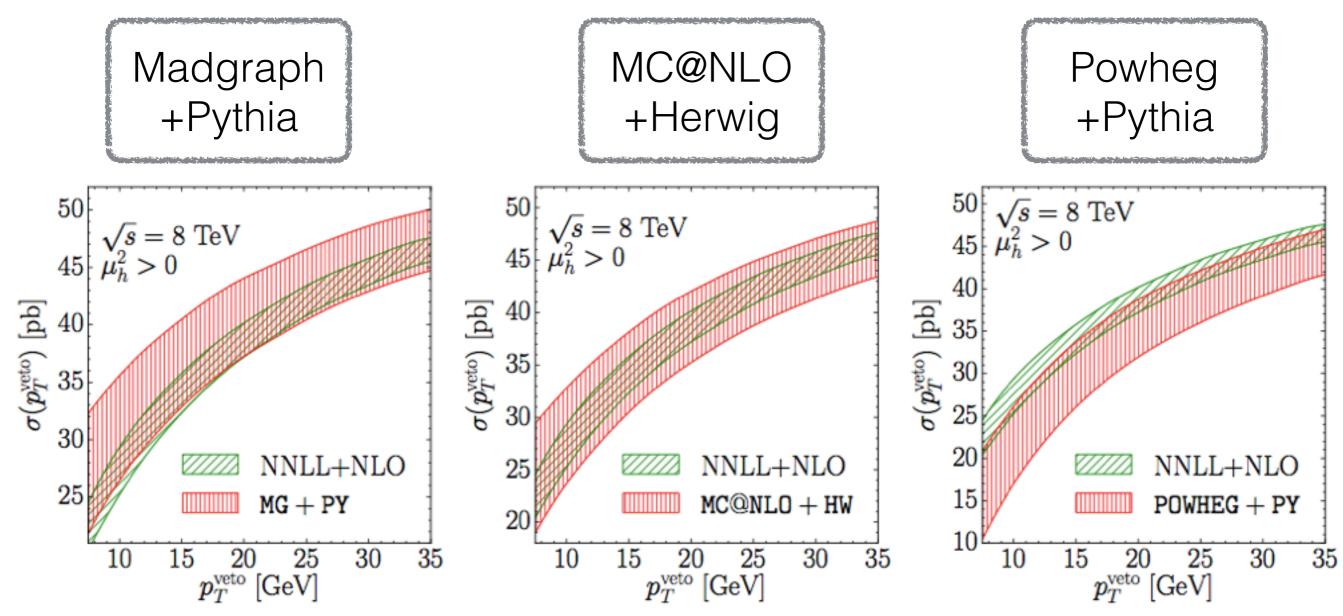
Splits the phase space integrals into regions of different rapidities

- Corresponding to rapidity regulator exists a rapidity renormalization group (RRG) equations, just as there exists RG equations corresponding to dim reg regulator µ.
- Formulation of RRG with analytic regulator had been missing in the literature. We address this issue [arXiv:1506.07529: PJ and Takemichi Okui].



More robust scale uncertainties using RRG (central value unchanged).

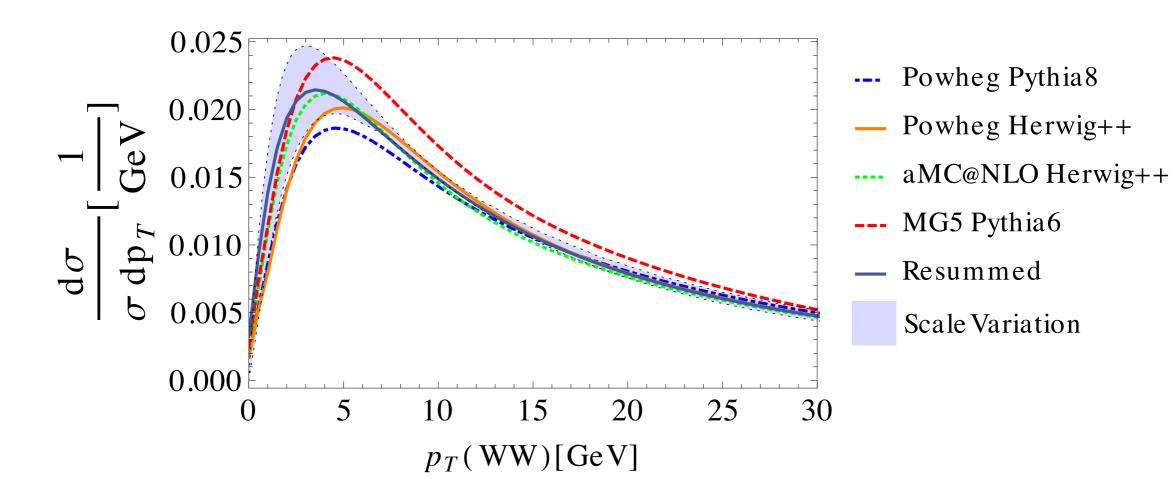
Comparison with MC+Parton shower



*without π^2 Resummation

Method II: p_T Resummation

 Discrepancy between p_T (WW) distribution shapes from NNLL p_T-resummation and MC [arXiv:1407.4481, P. Meade et al.]



Reweight MC :

$$F[\xi] = \frac{\text{Resummed bin}[\xi]}{\text{MC bin}[\xi]}$$

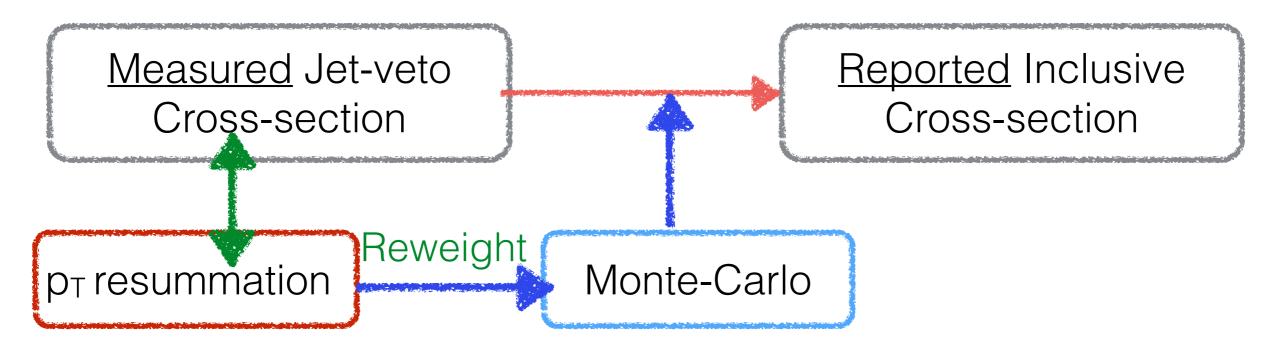
Method II: p_T Resummation

• **New CMS 8 TeV analysis** [CMS-PAS-SMP-14-016] reweights MC to correct for the p⊤ distribution.

$$\sigma_{W^+W^-} = 60.1 \pm 0.9 \text{ (stat.)} \pm 3.2 \text{ (exp.)} \pm 3.1 \text{ (th.)} \pm 1.6 \text{ (lum.) pb.}$$

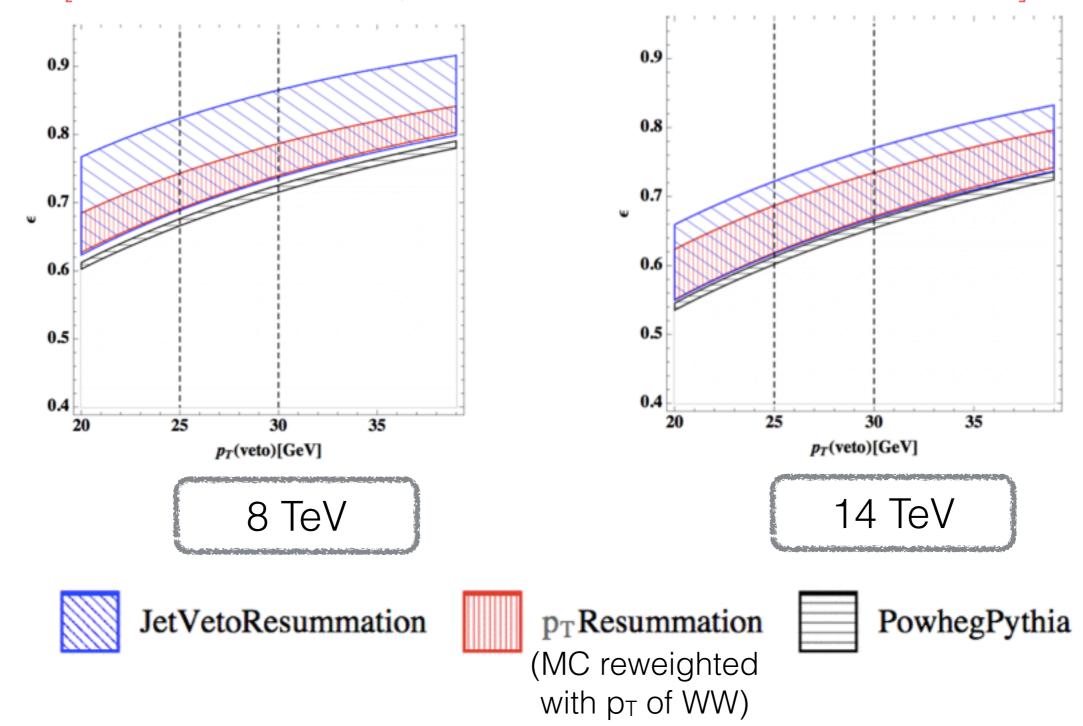
Theory:
$$59.8^{+1.3}_{-1.1} \, pb$$

 Some correlations between jet-veto and pT of the WW system captured by p_T reweighting technique.

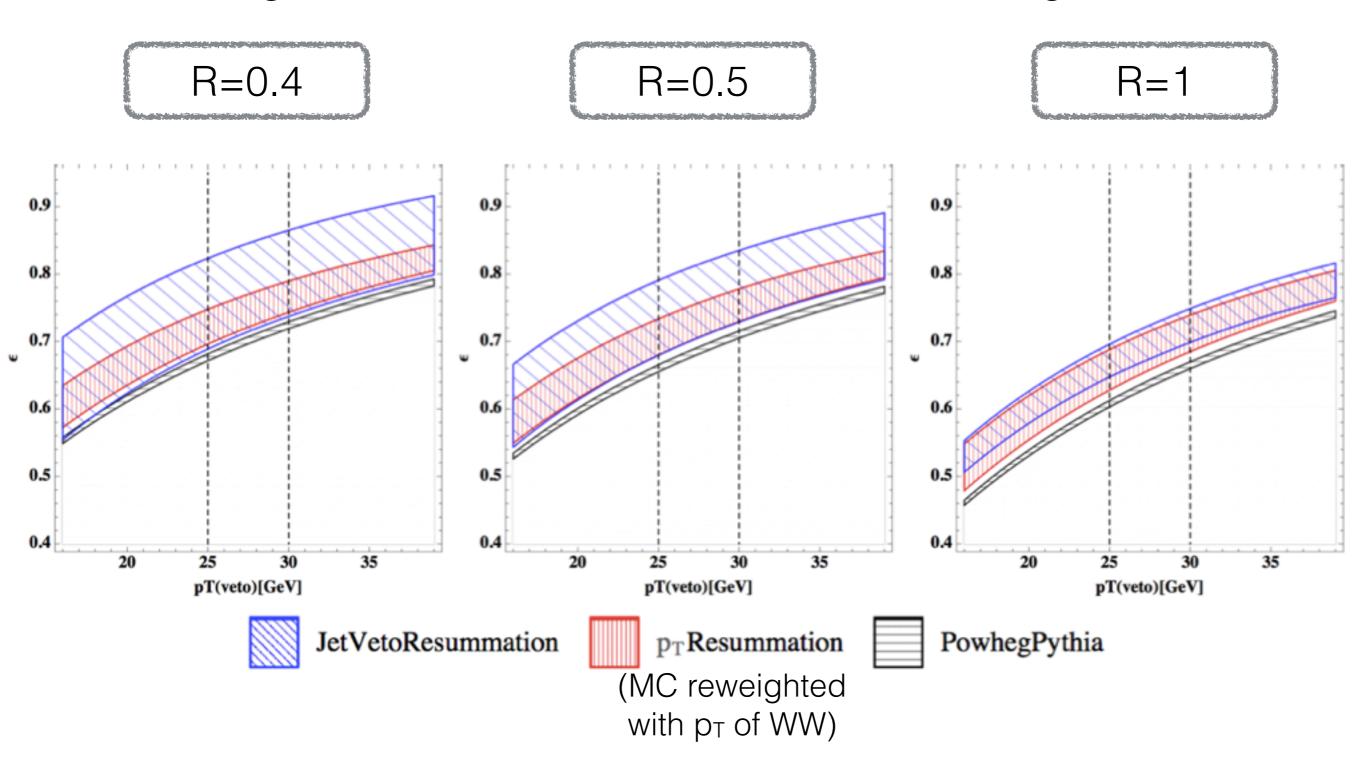


- Jet-veto resummation: Guaranteed to get the correct total jetveto cross-section but provides no information on the p_T shape.
- p_T resummation: Guaranteed to get the correct p_T shape for the total cross-section but relies on MC for jet-binning.
- Main source of discrepancy between two methods (~ 5-9%) is that <u>method I</u> employs π² resummation which also increases the inclusive cross-section while <u>method II</u> fixes the inclusive cross-section to fixed-order.
- Incidentally, increase in cross-section from π^2 resummation effect in <u>method I</u> is very similar to that from NNLO.

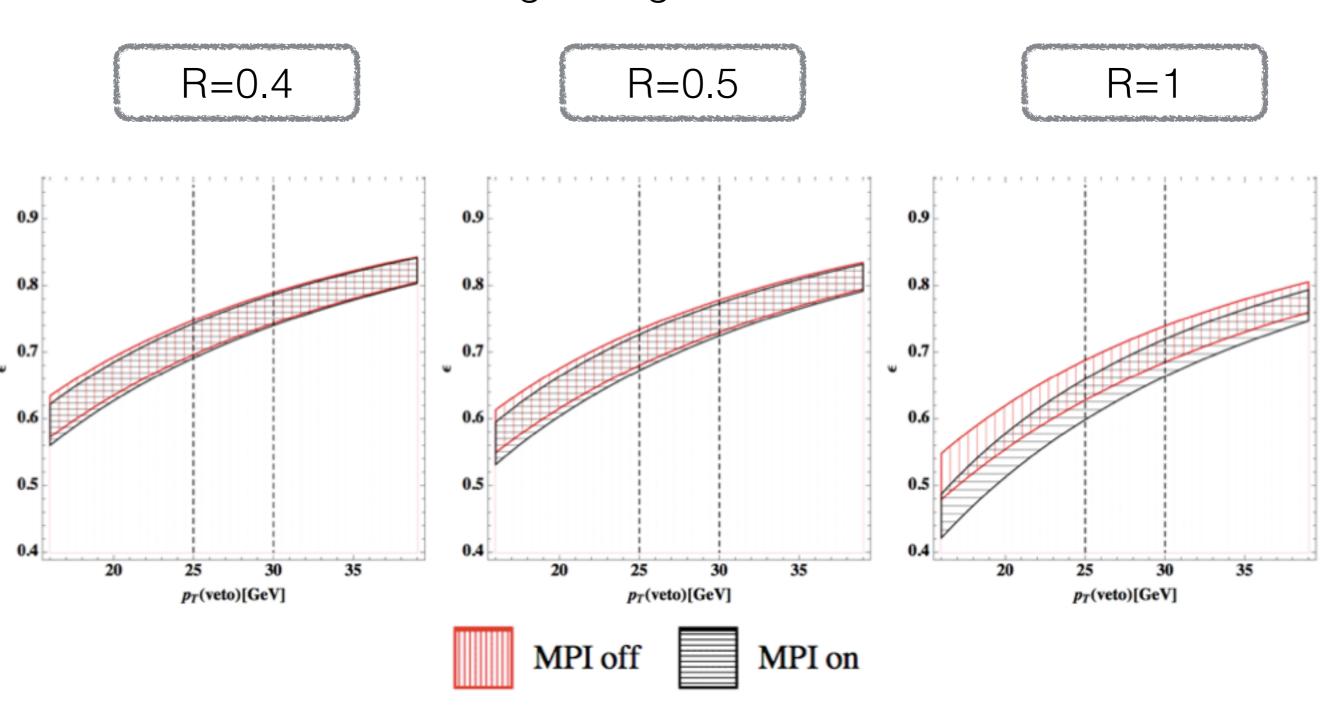
• For apples to apples comparison, use the same PDFs, scales, etc. [arXiv:1509.07118 : PJ, Patrick Meade and Harikrishnan Ramani]



Better agreement between the two methods for large R.



But MPI effects also big at large R.

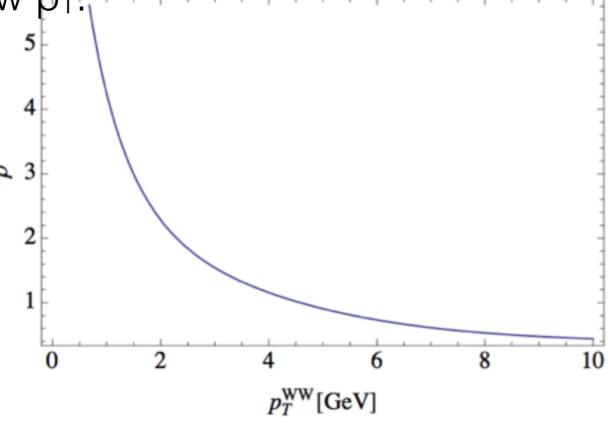


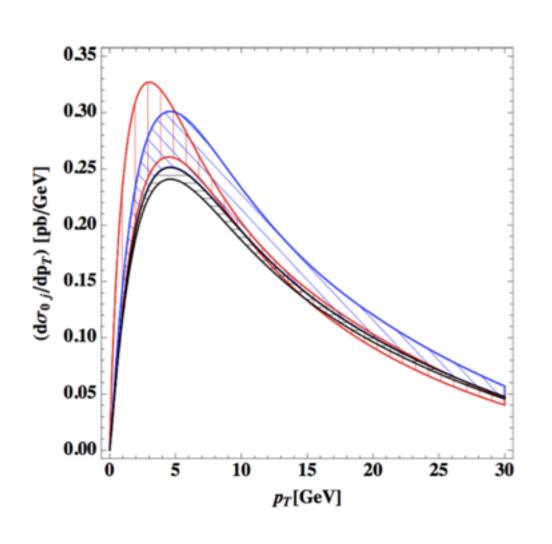
- Inspired by p_T -reweighting technique, we consider a new reweighting technique using of p_T of the leading jet instead of p_T of the WW system.
- The basic idea is to get the p_T shape of WW in the 0-jet bin using jet-veto resummation relying on the correlation between $p_T(WW)$ and $p_T(leading jet)$.

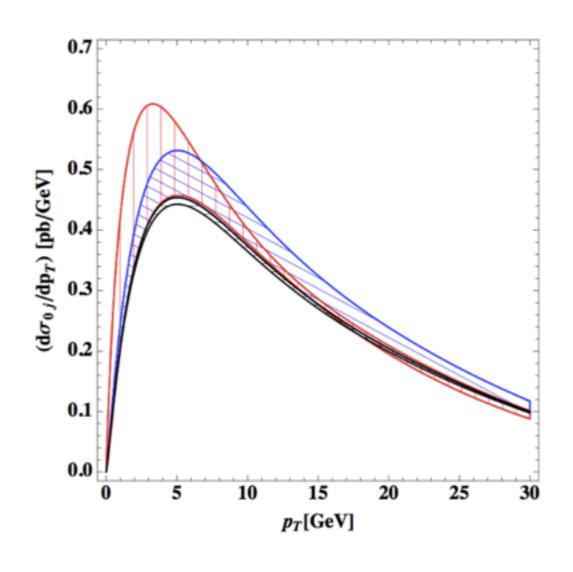
However, correlation worsens at low p_T.

$$\rho(p_T) = \frac{\langle |p_T^j(p_T) - p_T| \rangle}{p_T}$$

- Also, non-perturbative and MPI effects at low p_T lead to large uncertainties.
- Therefore, we simply use <u>two-bin</u> <u>reweighting</u> in p_T of the leading jet.







8 TeV

14 TeV

JetVetoResummation
(MC reweighted
with p_T of leading jet)



 p_T Resummation

(MC reweighted with p_T of WW)



PowhegPythia

Summary

- WW is an important background for Higgs studies and new physics searches.
- Jet veto in WW channel is often essential for suppressing top backgrounds.
- Jet veto logs can be large and need to be resummed.
- Two methods for resummation: Jet-veto resummation and p_T resummation followed by reweighting.
- Good agreement between two methods demonstrated and a new reweighting technique proposed.
- Analysis should be extended to other diboson channels for better understanding.